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# Quartz solubility in H<sub>2</sub>O-NaCl and H<sub>2</sub>O-CO<sub>2</sub> solutions at deep crust-upper mantle pressures and temperatures: 2–15 kbar and 500–900°C

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**ABSTRACT**—The solubility of quartz in  $H_2O$ -NaCl solutions was measured at 2, 4.35, 10 and 15 kbar and 500–900°C, and at NaCl concentrations up to halite saturation, usually greater than 75 wt.%. Quartz solubility was also measured in  $CO_2$ - $H_2O$  solutions at 10 kbar and 800°C. Solubilities were determined by weight loss of ground and polished quartz crystal fragments which were equilibrated with solutions in Pt envelopes for one to four days and then rapidly quenched. Experiments at 2 kbar were made with externally heated cold-seal apparatus; higher pressure experiments were done in a  $\frac{3}{4}$  inch-diameter piston-cylinder apparatus with NaCl pressure medium and graphite heater sleeve. Equilibrium solubility was demonstrated in several ways, and the present results reproduce those of Manning (1994) in pure  $H_2O$  at selected conditions.

At pressures below 4 kbar, NaCl in solution causes an initial "salting-in", or quartz solubility enhancement, which, at 2 kbar and 700°C, persists to concentrations as great as 70 wt.% NaCl before quartz solubility again becomes as low as in pure H<sub>2</sub>O. The maximum solubility occurs at  $X(H_2O) \sim 0.9$  and is 50% higher than in pure H<sub>2</sub>O. At 4.35 kbar and 700°C, however, quartz solubility decreases slightly with initial NaCl concentration, and then begins to drop rapidly with increasing salinity beyond 45 wt.% NaCl. At 10 and 15 kbar there is a steep initial decline in silica molality at all temperatures in the range 500–900°C, leveling off at higher NaCl concentrations. There is thus a pronounced change in solution behavior with pressure, from initial salting-in below 4 kbar to monotonic salting-out above 5 kbar. This pressure-induced change in silica solubility parallels the sharp decrease in H<sub>2</sub>O activity in NaCl solutions in the same pressure range found by Aranovich and Newton (1996). Therefore, the pressure-induced change in silica solubility is inferred to be a consequence of the dissociation of the neutral NaCl° complex to Na<sup>+</sup> and Cl<sup>-</sup> as solution densities increase above about 0.7 gm/cm<sup>3</sup>. At very high salinities, approaching halite saturation, the isobars of quartz solubility as a function of molten salts, pressure has only a minor effect on quartz solubility.

Quartz solubility at 10 kbar shows exponential decline with increasing salinity at all temperatures in the range 500°C to 900°C. This is the expected behavior of a two-component solvent, in which quartz is sparingly soluble in one component. At 10 kbar, isotherms of log silica molality versus  $H_2O$  mole fraction are linear between  $X(H_2O) = 1.0$  and 0.5, but begin to curve to lower values at 900°C, where high salinities are attained before halite saturation occurs. This behavior implies that the solute silica species is a hydrate that becomes progressively destabilized at low  $H_2O$  concentrations of the solvent. Plots of log silica molality versus  $\log H_2O$  activity suggest that the solute species is neutral  $H_4SiO_4$  with no additional solvated  $H_2O$  molecules, assuming no Na-SiO<sub>2</sub> complexing.

The solubility of quartz in  $CO_2$ -H<sub>2</sub>O fluids at 800°C and 10 kbar is much smaller than in NaCl solutions at the same P,T and H<sub>2</sub>O activity. Thermodynamic analysis suggests that the solute species in  $CO_2$ -H<sub>2</sub>O fluids is H<sub>4</sub>SiO<sub>4</sub> with 1–3 solvated H<sub>2</sub>O molecules, which is similar to the solute behavior inferred by Walther and Orville (1983) in  $CO_2$  and Ar solutions with H<sub>2</sub>O at lower pressures.

The present results show that  $SiO_2$  will partition very strongly into a concentrated salt solution in deep crust-upper mantle metamorphic and metasomatic processes, in preference to a coexisting immiscible  $CO_2$ -rich fluid. The much greater permeability of silicate rocks for salt solutions than for  $CO_2$ -rich solutions, together with the much higher solubility of silica-rich phases in the former, could be an important factor in geochemical segregation processes involving rising and cooling fluids of magmatic or metamorphic origin. *Copyright* © 2000 Elsevier Science Ltd

### 1. INTRODUCTION

The aqueous solubility of silica is one of the most important parameters in modeling metasomatism and fluid-rock interactions in the outer parts of the earth. Calculations of mass transport in regional metamorphism (Ague, 1994), oceanic hydrothermal systems (Lowell et al., 1993), and subduction complexes (Sorensen, 1988; Manning, 1995; 1997) require knowledge of silica concentration in migrating fluids under a broad range of pressure-temperature conditions.

The major source of information on quartz solubility as a function of pressure and temperature has been experimental studies in pure water, starting with the low-pressure studies of Kennedy (1950) and Morey and Hesselgesser (1951), through the work of Anderson and Burnham (1965) to 9 kbar and 900°C, to that of Manning (1994) to 20 kbar and 900°C. All of these studies used the weight-loss method, in which a ground fragment or crushed aggregate of natural quartz was equilibrated with initially pure water for several hours to a few days

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and then swiftly cooled to room temperature in several seconds to tens of minutes, depending on the type of apparatus. The combined dataset of earlier observations was summarized by Manning (1994) in a single equation, in which the solubility of quartz is expressed as an empirical function of temperature and water density.

Natural aqueous fluids are not pure water, but complex mixtures with other volatile species, such as CO2, and dissolved mineral components, especially alkali and alkaline earth halides and carbonates. Sampling of fluids in geothermal areas (e.g., Fournier, 1987) and modern fluid inclusion studies have shown that NaCl-dominated brines occur in many different petrogenetic settings. The latter technique has revealed that concentrated brines are associated with some sulfide ore bodies (Roedder, 1971), alkaline igneous rocks (Lowenstern, 1994; Belkin et al., 1997), carbonatites (Samson et al., 1995), highgrade metamorphic rocks (Smit and Van Reenen, 1997), peridotites, as at Zabargad Island, Red Sea (Sciuto and Ottonello, 1995), alpine eclogites (Philippot and Selverstone, 1991), and the ultra-high-pressure metasediments of the Western Alps (Philippot et al., 1995). Saline fluid inclusions from these different kinds of rocks typically contain halite and/or sylvite daughter crystals and have bulk compositions with NaCl contents as high as 80 weight percent. These hyperconcentrated brines could equally well be termed "hydrosaline magmas" (Roedder, 1971). Suites of apparently coeval concentrated brines and CO2-rich fluid inclusions have been reported from some high-grade rocks (Crawford and Hollister, 1986; Touret, 1985).

Experiments on synthetic fluid inclusions in the system  $H_2O-CO_2$ -NaCl in quartz crystals indicate that immiscible brine and  $CO_2$ -rich fluids can coexist at pressures of 7 kbar and temperatures of at least 900°C (Gibert et al., 1998). The two contrasting types of fluid have greatly differing mechanical properties in relation to silicate rocks.  $CO_2$ -rich fluids have low grainwetting ability (high dihedral angles), whereas concentrated salt solutions have extraordinary infiltration ability (Watson and Brenan, 1987). Partitioning of silica and other major and trace components of rocks between the immiscible fluids is therefore an important consideration for modeling metasomatic processes.

The solution of quartz into a host fluid is described formally by the reaction:

$$SiO_2$$
 (quartz) +  $nH_2O = SiO_2 \cdot nH_2O$  (aqueous) (1)

The solute species is thought to be dominantly a neutral complex represented by the formulas  $Si(OH)_4$  or  $H_4SiO_4$  with one to three hydrogen-bonded (solvated) water molecules. Thus, n in Eqn. (1) may be as high as 5. The hydration state of aqueous silica was investigated by Walther and Orville (1983) from existing quartz solubility data in the system  $CO_2$ - $H_2O$  and their own data for Ar- $H_2O$ . Most of the experimental data in  $CO_2$ - $H_2O$  solutions are contained in an unpublished M.S. thesis (Shettel, 1974), which presents experimental data to 5 kbar and 900°C. They reveal generally low solubility in  $CO_2$ -rich fluids, increasing with T and P at constant  $X(H_2O)$ . Most of the data plot with slopes of 3 to 4 on isothermal, isobaric diagrams of log silica molality versus log  $H_2O$  activity, which Walther and Orville (1983) interpret as indicating that the dominant dissolved molecular silica species has one to two solvated  $H_2O$  molecules. This analysis assumes that there is no solute interaction (complexing) between dissolved silica and the diluting volatile,  $CO_2$  or Ar.

Two previous studies are available on the solubility of quartz in concentrated NaCl solutions at supercritical temperatures and pressures. The Novgorodov (1977) measurements at 1.5 kbar and 700°C cover the salinity range 0-13 mole% NaCl. He found that NaCl substantially increases the solubility of quartz at these conditions. This unique "salting in" is most pronounced at low NaCl concentrations. Xie and Walther (1993) present measurements in dilute NaCl solutions (maximum of 0.83 molal or 1.5 mole% NaCl). The maximum temperature and pressure of the study are 2 kbar and 573°C. They found that NaCl initially increases the solubility of SiO<sub>2</sub> over that in pure  $H_2O$  at 1 kbar at >300°C and decreases it at <300°C. This behavior is quite in contrast to that of CO<sub>2</sub> solutions. Anderson and Burnham (1967) found a similar solubility enhancement in KCl solutions at 3 kbar and 600°C. Such salting-in behavior is observed to a more limited extent at 150-350°C at or near H<sub>2</sub>O liquid-vapor saturation (e.g., Chen and Marshall, 1982; Fournier et al., 1982; Saccocia and Seyfried, 1990).

Two distinct mechanisms were proposed to explain the silica "salting-in" effect. Anderson and Burnham (1983) invoked a reaction between solute NaCl or KCl and quartz to produce an alkali-silica hydrous complex and hydrochloric acid:

$$SiO_2 (quartz) + NaCl + H_2O = H_3NaSiO_4 + HCl$$
(2)

No salting-in effect exists at 25°C and one bar; therefore Anderson and Burnham (1983) concluded that increasing temperature and/or pressure promote salting-in. Xie and Walther (1993) ascribed the salting-in to the action of charged ions on the dielectric constant of the solution, without complex formation other than solvation. Their experiments suggest that solution density, rather than temperature and pressure independently, may be a controlling factor, in that salting-in begins at higher temperature at higher pressure: the change from saltingout to salting-in occurs near 300°C at 0.5 kbar, and above 400°C at one kbar. The cross-over density may therefore be about 0.7 gm/cm<sup>3</sup>.

The present study extends the measurements of silica solubility in NaCl solutions to the temperature and pressure regimes of the deeper crust and upper mantle, and to the higher salinity ranges indicated by the fluid inclusion studies on rocks of deeper provenance. Particularly important is the characterization of the effects of temperature, pressure and solution composition on the silica solution enhancement. Another important correlation to be made is with the pronounced pressure-induced decrease of H<sub>2</sub>O activity in NaCl and KCl solutions above 3-4 kbar found by Aranovich and Newton (1996, 1997). This activity change could affect the solution properties of SiO<sub>2</sub> by either the mechanism proposed by Anderson and Burnham (1983) or that of Xie and Walther (1993). We have measured quartz solubility in NaCl solutions at conditions of 2-15 kbar and 500-900°C and NaCl concentrations up to halite saturation (at most conditions, greater than 75 wt.% NaCl). We also determined quartz solubility in CO<sub>2</sub>-H<sub>2</sub>O solutions at 10 kbar and 800°C by the same method to compare quartz-saturated silica contents and solution mechanisms in the contrasting fluids.

#### 2. EXPERIMENTAL METHODS

The methods of the present study follow those of Manning (1994), except for a substantial reduction of the fluid volumes and crystal weights. Experiments from 4 to 15 kbar were made in a 3/4 inch diameter piston-cylinder apparatus with NaCl pressure medium and cylindrical graphite heater sleeve. Experiments at 10 and 15 kbar were made under piston-out conditions; that is, assemblies were pressed cold to a pressure 3 kbar less than the desired final pressure. Thermal expansion of the assemblies in heating increased the pressure to the final value. During heating at temperatures above 600°C, some bleeding of pressure was required to maintain the desired value. Under these conditions the pressure apparatus is virtually frictionless; that is, no pressure correction need be applied to the gauge or nominal pressure (Johannes, 1973; Mirwald et al., 1975). In all experiments, it was necessary to consolidate the assemblies at a minimum of 5 kbar before heating so that a large amount of bleeding was required to reduce the pressure to a nominal 4 kbar at the final temperature. The highest sample temperature readily obtainable at 4 kbar was limited to 750°C by the extensional stresses produced in the decompression.

The anomalous pressure-temperature path necessary to achieve a pressure below 5 kbar in the piston-cylinder apparatus makes independent pressure calibration necessary. Calibration was based on the P-T curve for the brucite to periclase dehydration reaction, as accurately determined in a gas-pressure apparatus by Aranovich and Newton (1996). Capsules containing coarsely crystalline periclase made by an arc-fusion process plus excess H2O exactly emulated the experimental capsules of the present study. Calibration experiments were of the same duration as the solubility measurements: one to two days. If a quenched capsule contained abundant coarse brucite, the stability field of that mineral was considered to have been encountered, inasmuch as the transition is quite rapid with small temperature overstep above 600°C and several kbar. Calibration experiments at a nominal pressure of 4.0 kbar yielded an equilibrium reaction temperature of 727  $\pm$  3°C, corresponding to a pressure of  $4.35 \pm 0.15$  kbar. Maximum pressure uncertainties at all pressures are believed no greater than  $\pm 0.3$  kbar.

Temperature was measured and controlled by calibrated matched pairs of W-3%Re versus W-25%Re thermocouples. The stated accuracy of the manufacturer (Engelhard) was better than 1°C. A precise digital millivoltmeter in parallel with a Love proportional controller consistently showed a temperature ripple of  $\pm$ 1°C, with a period of 30 sec. The thermocouple was in virtual contact with the quartz crystal: typically, the thermocouple ip made a slight dent in the Pt sample envelope. Maximum uncertainty in sample temperatures from all sources was  $\pm$ 3°C. Quench times were very fast: temperatures fell below 200°C in about 12 sec. after turning off the heating power.

Experiments at 2 kbar were made in externally heated cold-seal bombs. The bombs were mounted horizontally, with the samples at the pre-determined plateau of temperature. Temperatures were measured and controlled by external chromel-P-alumel thermocouples inserted in a well in the bomb at the sample position. Internal metal rods filled the bomb cavity adjacent to the samples. Pressures are accurate to  $\pm 50$  bars, and temperatures to  $\pm 3^{\circ}$ C. Quench was by air-blow, and required three minutes to bring the temperature to below 200°C. Pressure was maintained during the quench.

Two large, gemmy natural quartz crystals, one from Brazil (Manning, 1994) and the other from Mountain Valley, AR, provided samples. Small equant chips were ground into ellipsoidal shapes by rolling between sheets of 240-grit emery paper. The crystals were then rolled between sheets of 600-grit paper until they acquired a fair polish. Weights of experimental crystals were mostly 0.3 to 3.5 mg. Larger crystals occasionally broke during the experiments, sometimes during the quench or upon decompression, to judge from unetched conchoidal fractures. A weighed quartz ellipsoid was placed in a 1.3 cm long Pt tube segment of 2 mm diameter and 0.18 mm wall thickness with weighed amounts of multiply-distilled and deionized H<sub>2</sub>O, usually about 10 mg, and reagent-grade (Baker ACT) NaCl, previously dried for an hour at 310°C. The capsule was sealed by arc-welding and folded double to produce a packet 2 mm thick and 5 mm square. All weighings were made on a Mettler M3 microbalance, with  $1\sigma$  reproducibility of 0.002 mg. H<sub>2</sub>O contents were determined in two ways, by weighing after loading with a microliter syringe, making a calibrated allowance for loss by evaporation, and by weight loss of punctured quenched capsules dried for 15 min. at  $115^{\circ}$ C and then for 15 min. at  $310^{\circ}$ C. The two methods generally yielded the same H<sub>2</sub>O contents within 0.5%. The puncture-and-dry weighings were considered to be the more reliable, except for an occasional large discrepancy, which could usually be traced to loss of NaCl from a capsule during drying.

The experiments regarded as definitive were those in which a single unbroken quartz crystal was retrieved from the quenched capsule. Quartz solubility in these experiments was determined from the weight lost by the retrieved crystal. In a few runs, one or a number of small terminated quartz crystals nucleated and grew independently of the seed crystal, probably because of a small temperature gradient across the capsule. Inasmuch as it was usually impossible to be sure that all such small crystals were accounted for, the net weight loss was uncertain. These runs were repeated, as were the few runs in which there was a detectable weight loss of the capsule, signifying a leak.

Experiments at the lowest  $H_2O$  percentages were halite-saturated, as shown by the presence of a single large rounded halite crystal, in addition to the quartz crystal, amid the mass of small skeletal salt crystals quenched from the fluid phase. Halite saturation limits the NaCl concentration at a fixed pressure and temperature, but affords the advantage that the fluid compositions at halite saturation are well known, as determined experimentally by Koster Van Groos (1991) and Aranovich and Newton (1996).

A number of experiments were made with quartz crystals sealed with mixtures of liquid water, oxalic acid dihydrate, 2COOH.2H2O, and silver oxalate, Ag<sub>2</sub>C<sub>2</sub>O<sub>4</sub>, to generate fluids on the join H<sub>2</sub>O-CO<sub>2</sub>. Under conditions of high oxygen fugacity, oxalic acid dihydrate yields a fluid of equimolar CO2 and H2O, with outward diffusion of hydrogen through the Pt capsule walls. Silver oxalate yields pure CO<sub>2</sub> and silver metal. Fluid ingredients plus a quartz crystal were sealed without significant weight loss in short segments of 1-mm-diameter Pt tubing. Because of space limitations, crystal weights were only 0.3-0.5 mg. Typical volatile contents were 3-5 mg. The Pt capsules were in turn loaded into segments of 3 mm diameter Au tube with about 40 mg of hematite and 4 mg of H2O, and sealed by welding. The resulting buffer capsules had geometries nearly identical to the larger Pt capsules of the unbuffered runs. At the conclusion of the runs, the quenched buffer capsules were checked for leakage, and to verify that hematite and magnetite were both present, constituting an HM buffer. Retrieved Pt capsules were inflated with CO<sub>2</sub> pressure, unless there was leakage during or after the run. Fluids were analyzed by the puncture-weight loss technique described by Aranovich and Newton (1999). First, a capsule was frozen in liquid nitrogen and punctured with a needle while still frozen. During warming, an immediate weight loss was ascribed to CO2 escape. The capsules were dried at 115°C and then at 310°C for 15 min. each and reweighed to obtain the H2O content. The H2O content determined in this way agreed with the initially loaded value to within  $\pm 0.12$  mg.

#### 3. RESULTS OF SOLUBILITY MEASUREMENTS

# **3.1.** Weighing Uncertainties, Reproducibility, and Equilibration

Tables 1 and 2 list data for the experiments with NaCl- and CO<sub>2</sub>-bearing aqueous solutions, respectively, including molality values calculated from (crystal weight loss)\*1000/ (60.085\*(weight H<sub>2</sub>O)). Fluid inclusions trapped during experiments were too few and tiny to influence solubility determinations. Uncertainties due to weighing were estimated numerically. Replicate weighings of solids yielded a normal distribution of weighing errors with  $1\sigma = 0.002$  mg. The difference between the initial H<sub>2</sub>O weight, determined by applying an evaporation correction, and the weight determined after each experiment by drying was also normally distributed and averaged -0.065 mg with  $1\sigma = 0.079$  mg. This discrepancy suggests that we overcorrected for evaporation loss, either because addition of NaCl after water decreased the rate of evaporation from the capsule or, despite drying, a small amount

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Table 1. Quartz solubility measurements in H<sub>2</sub>O-NaCl fluids.

Expt. no.	P (kbar)	T (°C)	Time (hr)	Initial qz wt. (mg)	Final qz wt. (mg)	Initial H <sub>2</sub> O wt. (mg)	Initial NaCl wt. (mg)	$X_{H_2O}$	a <sub>H2</sub> O	X <sub>NaCl</sub>	10 <sup>3</sup> *X <sub>SiO2</sub>	m <sub>SiO2</sub> (mol/kgH <sub>2</sub> O)
45	2	700	91	0.306	0.247	7.057	0.000	0.997	0.997	0.000	2.5 (3)	0.139 (15)
56	2	700	96	1.713	1.567	11.301	2.449	0.934 (1)	0.944 (1)	0.062(1)	3.6 (2)	0.215 (10)
48	2	700	96	2.150	2.057	7.403	4.497	0.840 (4)	0.862 (4)	0.157 (4)	3.2 (2)	0.209 (15)
47	2	700	96	0.483	0.402	9.307	17.589	0.631 (6)	0.671 (6)	0.368 (6)	1.6(1)	0.145 (11)
49	2	700	96	2.792	2.759	7.036	24.377	0.483 (8)	0.527 (8)	0.516 (8)	0.7 (1)	0.078 (15)
61	4.35	600	46	0.873	0.725	11.276	0.000	0.996	0.996	0.000	3.9 (2)	0.219 (10)
59	4.35	600	43	2.019	1.929	11.108	14.531	0.711 (4)	0.594 (4)	0.287 (4)	1.7 (1)	0.135 (9)
56	4.35	600	47	1.783	1.722	10.741	20.332	0.631 (5)	0.504 (5)	0.368 (5)	1.1 (1)	0.095 (9)
43	4.35	700	19	0.927	0.714	9.421	0.000	0.993	0.993	0.000	6.7 (3)	0.377 (14)
60	4.35	700	21	1.719	1.477	11.066	3.345	0.909 (2)	0.884 (2)	0.085 (2)	6.0 (2)	0.364 (12)
46	4.35	700	23	1.851	1./10	6.993	4.424	0.833(4)	0.789 (4)	0.162(4)	5.0(3)	0.336 (18)
/0	4.35	700	23	1.608	1.431	10.836	11.721	0.747(4)	0.688(4)	0.249(4)	3.7(1)	0.272(11)
41	4.35	700	18	1.932	1.849	6.809	11.295	0.660(7)	0.591(7)	0.338(7)	2.4(2)	0.203(16)
42	4.35	700	12	0.208	0.233	4.932	10.185	0.497 (11)	0.424 ((11)	0.502 (11)	1.1(2)	0.118(21) 0.458(0)
72 60	4.55	750	13	2.173	1.720	18.200	18,000	0.992 0.762 (2)	0.992 0.726 (2)	0.000	6.2 (2) 5 1 (1)	0.438(9) 0.268(7)
09 77	4.55	750	22	1.929	1.520	7 160	34 158	0.703(2) 0.405(7)	0.720(2) 0.358(7)	0.232(2) 0.594(8)	3.1(1)	0.308(7) 0.112(15)
M9/	10	500	25	10.0/15	10 201	63 688	0.000	0.403(7)	0.338(7)	0.000	3.1	0.112(13) 0.171(2)
M94	$\frac{10}{10}$	500	$2\frac{23}{62}$	9 4 5 2	8 684	73 952	0.000	0.997	0.997	0.000	$\frac{3.1}{3.1}$	$\frac{0.171(2)}{0.173(2)}$
68	$\frac{10}{10}$	500	$\frac{202}{70}$	$\frac{7.432}{0.726}$	$\frac{0.004}{0.620}$	$\frac{13.952}{14.795}$	<u>6.061</u>	$\frac{0.997}{0.886}(2)$	$\frac{0.757}{0.759}$ (2)	$\frac{0.000}{0.112}(1)$	$\frac{3.1}{1.9}(1)$	$\frac{0.173(2)}{0.119(7)}$
63	10	500	93	0.763	0.020	10 783	16 174	0.000(2)	0.759(2)	0.22(2)	1.7(1)	0.077(9)
67	10	600	72	1.083	0.727	16.817	0.000	0.994	0.994	0.000	6.3(1)	0.353 (8)
37	10	600	69	1.013	0.880	10.995	6.948	0.834 (3)	0.688 (3)	0.163 (3)	3.0(1)	0.201 (10)
40	10	600	48	0.293	0.209	11.148	19.372	0.67(2)	0.47(2)	0.33(2)	1.5(1)	0.125 (9)
52	10	700	24	2.068	1.608	11.201	0.000	0.988	0.988	0.000	12.2 (3)	0.684 (17)
62	10	700	22	1.154	0.797	10.992	2.825	0.918(1)	0.850(1)	0.073 (1)	8.9 (2)	0.541 (15)
31	10	700	44	1.037	0.777	10.776	6.695	0.834 (3)	0.710 (3)	0.160 (3)	6.0 (2)	0.402 (13)
30	10	700	24	1.323	1.154	10.513	15.067	0.691 (5)	0.518 (5)	0.305 (4)	3.3 (1)	0.268 (11)
32	10	700	72	0.618	0.559	6.238	24.384	0.56(2)	0.38(2)	0.44 (2)	1.6 (2)	0.158 (17)
20	10	800	23	2.108	0.446	22.176	0.000	0.978	0.978	0.000	22.0 (2)	1.248 (13)
26	10	800	20	1.971	1.310	11.084	2.408	0.922 (1)	0.874 (1)	0.062(1)	16.5 (3)	0.993 (22)
14	10	800	28	2.035	0.830	22.099	8.451	0.882 (1)	0.808 (1)	0.104 (1)	14.4 (1)	0.908 (10)
25	10	800	24	0.800	0.292	10.409	5.705	0.845 (3)	0.751 (3)	0.143 (3)	12.4 (3)	0.813 (20)
9	10	800	26	2.849	2.424	10.848	8.900	0.791 (3)	0.673 (3)	0.200 (3)	9.3 (2)	0.652 (17)
19	10	800	24	1.274	0.907	10.787	11.698	0.744 (4)	0.610 (4)	0.249 (4)	7.6 (2)	0.567 (15)
11	10	800	72	3.618	3.298	11.095	14.910	0.703 (4)	0.560 (4)	0.291 (4)	6.1 (1)	0.480 (14)
16	10	800	23	0.593	0.260	11.149	15.732	0.693 (4)	0.548 (4)	0.301 (4)	6.2 (1)	0.497 (14)
27	10	800	20	1.3/4	1.082	10.164	14.542	0.690 (5)	0.544 (5)	0.304(4)	5.9 (2)	0.478 (15)
18	10	800	46	2.746	2.470	11.212	19.962	0.643(5)	0.491 (5)	0.353(4)	4.7(1)	0.410(12)
24	10	800	27	1.442	1.232	10.502	24.281	0.582(5)	0.427(5) 0.201(5)	0.415(5)	3.5(1)	0.333(13) 0.202(11)
17	10	800	22	0.901	0.705	10.807	29.108	0.340(3)	0.391(3)	0.431(3)	3.0(1)	0.305(11) 0.270(11)
21	10	800	23 53	2.323	2.330	11.149	34.734 41.570	0.309(3)	0.330(3)	0.469(3) 0.531(5)	2.0(1)	0.279(11) 0.242(10)
15	10	800	<i>1</i> 6	1.024	0.540	11.272	54 282	0.407(3)	0.318(3) 0.28(2)	0.531(3)	2.0(1) 15(1)	0.242(10) 0.202(10)
53	10	850	22	2 614	1 675	9 278	0.000	0.971(1)	0.20(2)	0.000	295(7)	1.685(41)
35	10	850	22	1 340	0.634	11 123	6 594	0.971(1) 0.832(3)	0.771(1) 0.744(3)	0.000 0.152(2)	15.8(3)	1.005(41) 1.057(23)
34	10	850	22	1.040	0.688	10 775	15 432	0.632(3) 0.688(4)	0.744(3) 0.557(4)	0.152(2) 0.304(4)	80(2)	0.648(17)
39	10	850	18	1.525	1.281	10.582	30.293	0.529(5)	0.389(5)	0.367(1) 0.467(5)	3.7(1)	0.384(13)
75	10	850	22	1.323	1.238	5.991	29.551	0.396(9)	0.270 (9)	0.602 (9)	1.7(1)	0.236 (19)
36	10	850	21	0.542	0.483	4.763	29.170	0.346 (11)	0.230(11)	0.653 (11)	1.3 (1)	0.206 (23)
74	10	875	20	1.727	1.674	4.773	48.170	0.30(2)	0.19(2)	0.70(2)	1.1 (1)	0.185 (23)
M94	10	900	17	13.85	7.92	41.25	0.000	0.959	0.959	0.000	41.3 (2)	2.39(1)
55	10	900	11	3.121	2.493	11.193	12.952	$\overline{0.728}$ (4)	$\overline{0.620}$ (4)	$\overline{0.260}$ (4)	12.3 (2)	$\overline{0.934(21)}$
54	10	900	10	2.360	2.172	7.513	26.699	0.476 (7)	0.352 (7)	0.521 (7)	3.6(1)	0.417 (19)
76	10	900	20	1.675	1.586	5.902	35.827	0.348 (8)	0.242 (8)	0.651 (8)	1.6(1)	0.251 (19)
64	10	900	14	3.039	3.012	2.947	44.556	0.24 (2)	0.16(2)	0.76(2)	0.7(2)	0.153 (36)
65	10	900	20	3.014	2.990	2.863	39.845	0.24 (2)	0.16(2)	0.76 (2)	0.6 (2)	0.140 (38)
<u>M94</u>	15	700	<u>45</u>	20.99	19.21	35.16	0.000	0.985	0.985	0.000	15.0(1)	0.843 (6)
79	15	700	20	1.190	0.807	10.946	3.435	0.903 (2)	0.797 (2)	0.087 (2)	9.5 (2)	0.583 (15)
50	15	700	22	1.248	0.982	9.496	5.490	0.843 (3)	0.685 (3)	0.150 (3)	7.1 (2)	0.466 (16)
78	15	700	20	1.588	1.406	10.271	13.784	0.705 (5)	0.488 (5)	0.292 (4)	3.7 (1)	0.295 (12)
51	15	700	25	1.576	1.482	6.724	12.595	0.632 (8)	0.406 (8)	0.365 (8)	2.7 (2)	0.233 (17)

Underlined entries from Manning (1994). Italics denote halite saturation. In halite-undersaturated experiments,  $X_i = n_i/(n_{H_2O} + n_{NaCl} + n_{SiO2})$ , whereas at halite saturation  $X_{H_2O}$  and  $X_{NaCl}$  were taken from Aranovich and Newton (1996) and ignore SiO<sub>2</sub> in the fluid.  $a_{H_2O}$  from Aranovich and Newton (1996) assuming that SiO<sub>2</sub> is an inert diluent. Parenthetical numbers are numerically propagated  $2\sigma$  weighing uncertainties in last digit(s) where >0.5 (see text) for all values except those at halite saturation, which also include an assumed  $2\sigma$  uncertainty in  $X_{H_2O}$  of 0.2.

Table 2. Quartz solubility measurements in H<sub>2</sub>O-CO<sub>2</sub> fluids.

Expt. no.	P (kbar)	T (°C)	Time (hr)	Initial qz wt. (mg)	Final qz wt. (mg)	Final H <sub>2</sub> O wt. (mg)	Final CO <sub>2</sub> wt. (mg)	$X_{H_2O}$	a <sub>H2O</sub>	X <sub>CO2</sub>	10 <sup>3</sup> *X <sub>SiO2</sub>	m <sub>SiO2</sub> (mol/kgH <sub>2</sub> O)
C-5	10	800	46	0.401	0.316	2.382	1.335	0.806(1)	0.863 (1)	0.185(1)	8.6 (6)	0.594 (40)
C-4	10	800	46	0.456	0.413	2.108	2.631	0.659(1)	0.771(1)	0.337(1)	4.0 (5)	0.340 (44)
C-1	10	800	23	0.414	0.391	1.835	4.482	0.499(1)	0.658(1)	0.499 (1)	1.9 (6)	0.209 (52)
C-6	10	800	45	0.295	0.278	1.740	4.322	0.495(1)	0.656(1)	0.503(1)	1.5 (5)	0.163 (55)
C-3	10	800	30	0.393	0.389	0.595	3.033	0.324 (2)	0.492 (1)	0.676 (2)	0.7 (9)	0.112 (156)

 $X_i = n_i/(n_{H_2O} + n_{CO_2} + n_{SiO_2})$ .  $a_{H_2O}$  from Aranovich and Newton (1999) assuming that SiO<sub>2</sub> is an inert diluent. Parenthetical numbers are numerically propagated  $2\sigma$  weighing uncertainties in last digit(s) (see text).

of H<sub>2</sub>O was retained in the amorphous silica quench-precipitate. All weighing-related errors were propagated to final concentrations by using a Monte Carlo routine in which 1000 trial weights were generated randomly using population characteristics determined from the weighing checks. Increasing the number of trial weights did not change the results. Final errors are reported in Tables 1 and 2 as  $2\sigma$  uncertainties that include weighing and evaporative loss of H<sub>2</sub>O during capsule loading. The uncertainty given in this study are less than that of Manning (1994), in which  $2\sigma = 0.040$  molal, because errors in that study were determined by pooling results of replicate experiments and therefore included uncertainties in pressure and temperature as well.

Two kinds of experiments were made which reproduced the Manning (1994) SiO<sub>2</sub> molality value at 10 kbar and 800°C in pure H<sub>2</sub>O. A single weight loss measurement (#20, Table 1) gave m(SiO<sub>2</sub>) =  $1.247 \pm 0.013$ , compared to the Manning (1994) value of  $1.232 \pm 0.040$  molal, demonstrating excellent agreement. The second kind of experiment introduced finely ground quartz into NaCl-H<sub>2</sub>O mixtures in progressively greater amounts in successive runs, with microscopic search for quartz crystals in the quenched charges. The onset of quartz saturation was signaled by the sudden appearance of a few large terminated quartz crystals. This method, though not as precise as the weight loss method, confirmed the validity of the latter method by yielding a molality of  $1.23 \pm 0.11$  in pure H<sub>2</sub>O at 800°C and 10 kbar.

Equilibrium solubility was indicated in two ways. A time study was made at  $X(H_2O)$  approximately 0.7 (runs 11, 16 and 27, Table 1). The initial weights of the quartz crystals ranged from 3.6 to 0.6 mg. Times ranged from 20 to 72 h. The results were identical within analytical errors. A more comprehensive proof of the attainment of equilibrium solubility was the appearance, in the majority of experiments, of broad, shiny crystal faces usually on one or both ends of the ellipsoidal quartz crystal, despite overall weight loss. This evidence of solution and reprecipitation is virtual proof of quartz saturation.

The two-kbar quenched crystals also showed incipient wellformed crystal faces. However, because of the much slower quench rates in the cold-seal apparatus, it is possible that there was some quench reprecipitation of dissolved silica on the crystals, which would result in underestimation of the solubility. Our measured silica molality in pure H<sub>2</sub>O is 0.139 at 700°C, compared to the Anderson and Burnham (1965) value of 0.158 molal at the same conditions (average of two runs).

### 3.2. Pressure Effect on Solubility and Salting-In

Figure 1 shows the solubility data at 700°C and 1–15 kbar. The drastic change in solubility behavior between 2 and 10–15 kbar is apparent. At 2 kbar there is a strong solubility enhancement effect, with quartz solubility rising steeply with NaCl concentration to a value about 50% higher at  $X(H_2O) = 0.9$  than in pure H<sub>2</sub>O. The initial rate of salting-in is about the same as at 1 kbar, if a small extrapolation of the data of Xie and Walther (1993) is made of their measurements at m(NaCl) = 0.49 and 0.83 (Fig. 1). The Novgorodov (1977) measurements show the same initial salting-in effect. The apparent crossing of his highest-salinity point at 1.5 kbar above our 2 kbar data could be an artifact of different experimental techniques, though an actual cross-over could conceivably occur in the



Fig. 1. Comparison of measurements of quartz solubilty in NaCl solutions at 700°C, 1–15 kbar, expressed as molality,  $m(SiO_2)$ , versus mole fraction of NaCl ( $X_{NaCl}$ ). The Xie and Walther (1993) points at 1 kbar were extrapolated from the data in their Fig. 2b. Solubility isobars are fit by eye. The position of halite saturation is uncertain because this condition was investigated only at 10 kbar at 700°C. Uncertainties that are smaller than symbol size are not plotted.



Fig. 2. Quartz solubility measurements at 10 kbar and 500–900°C. Data point of Manning (1994) at 900°C is off-scale. Except at 850°C, all data on the halite-saturation curve represent experiments in which a large halite crystal was present in the quenched charge.  $X_{\rm NaCl}$  in these halite-saturated experiments is from Aranovich and Newton (1996). Solubility isotherms are exponential fits to data at  $X_{\rm NaCl} < 0.6$  (see text and Fig. 3).

low-pressure region. At higher salinities, salting-out takes over, but the SiO<sub>2</sub> molality at  $X(H_2O) = 0.6$  is still about the same as in pure H<sub>2</sub>O. With further increase in salinity, approaching halite saturation, m(SiO<sub>2</sub>) is about half of the pure water value. Though the solubility values may be slightly affected by the slow quench in the cold-seal apparatus, the initial salting-in and solubility maximum are undoubtedly real. At 10 and 15 kbar there is only an exponential decrease of solubility with NaCl content of the fluid. At 4.35 kbar, the behavior is intermediate between these extremes, with an initially flat positive  $dm(SiO_2)/dX(H_2O)$  trajectory, turning with higher salinity into an exponential decline similar to the 10 kbar behavior. The quartz solubility enhancement effect found by Novgorodov (1977) and Xie and Walther (1993) is thus confined to low pressures and low NaCl concentrations. The large spread of solubility in pure H<sub>2</sub>O over 2–15 kbar converges to a narrow span in the high NaCl concentration range, ending in halite saturation.

#### 3.3. Exponential Salinity Dependence at High Pressure

At 10 and 15 kbar the solubility shows simple monotonic decrease with NaCl mole fraction at the temperatures investigated. Figure 2 shows the 10 kbar data. The decrease is exponential except possibly at the highest T and X(NaCl) studied. This behavior is clearly shown in the plot of Figure 3 of log  $m(SiO_2)$  versus X(H<sub>2</sub>O) at 10 kbar. The linearity is striking, especially in the near parallelism of the slopes at all temperatures. This regularity of the solubility implies that, at pressures



Fig. 3. Data of Fig. 2 on log m(SiO<sub>2</sub>) scale, demonstrating exponential dependence of m(SiO<sub>2</sub>) on H<sub>2</sub>O mole fraction, and that dlog m(SiO<sub>2</sub>)/dX<sub>NaCl</sub> is nearly independent of temperature. This relationship may break down at 900°C and X<sub>NaCl</sub> > 0.6, possibly because of destabilization of hydrous solute silica species at very low H<sub>2</sub>O activity.

approaching 10 kbar, a relatively simple medium for SiO<sub>2</sub> solution exists, while at lower pressures, the solvent is more complex, with composition- and pressure-dependent mixtures of solute species. Departure from the linear log m(SiO<sub>2</sub>) versus X(H<sub>2</sub>O) trend may be present in the highest-salinity point at 900°C and 10 kbar, where salinity can exceed 50 mole percent NaCl before halite saturation occurs. The halite-saturation boundary bends accordingly as silica solubilities begin to drop at X(H<sub>2</sub>O) < 0.5. The indication that SiO<sub>2</sub> solubilities become markedly lower at very low H<sub>2</sub>O activities implies that the dominant solute silica species is a hydrate. The nature of the limiting slopes in Figures 2 and 3 at very high salinities is uncertain: if SiO<sub>2</sub> is insoluble in pure molten NaCl, the slope is undefined.

# **3.4.** Temperature Independence of Relative Molality at High Pressure

The pressure-induced change in SiO<sub>2</sub> solution is shown even more dramatically in Figure 4. In the plot of  $m(SiO_2)/m_o(SiO_2)$ at 10 kbar, where m<sub>o</sub> refers to quartz solubility in pure H<sub>2</sub>O, the data at all temperatures describe a single exponential decline with X(NaCl). The curve through the points is an exponential fit to all 10 kbar data. No systematic temperature dependence of m/m<sub>o</sub> is evident. The data at 15 kbar fall nearly on the 10 kbar trend, and are not shown. The corresponding curve for 4.35 kbar is more complex. The experimental data for this pressure in the range 600–750°C describe a trend, though some temperature dispersion seems to be present. At 2 kbar and 700°C, the trend is grossly different because of the salting-in effect.

The observation that relative silica concentration is independent of temperature at high pressures provides a basis for a



Fig. 4. SiO<sub>2</sub> solubility relative to that in pure H<sub>2</sub>O at the same P,T conditions,  $m_x(SiO_2)/m_o(SiO_2)$ . The plot shows the extreme effect of pressure on silica solubility, in that salting-in (NaCl enhancement of solubility) occurs at 2 kbar, whereas exponential salting out is observed at 10 kbar. Behavior is intermediate at 4.35 kbar.

single expression for estimating quartz solubility in  $H_2O$ -NaCl fluids at high pressures and temperatures. We fit polynomials to the 4.35, 10, and 15 kbar data and then used interpolation to estimate the variation in relative molality with pressure at constant temperature. This led to the following equation:

$$\frac{\mathbf{m}_{\text{SiO}_2}}{\mathbf{m}_{0,\text{SiO}_2}} = \sum_{j=0}^{3} \sum_{i=0}^{4} v_{ij} \mathbf{X}_i \mathbf{P}_j$$
(3)

The parameters  $\nu_{ij}$  are given in Table 3 and the predicted variation of m(SiO<sub>2</sub>)/m<sub>o</sub>(SiO<sub>2</sub>) is given in Figure 5A. Figure 5B illustrates the differences between relative molality predicted by Eqn. (3) and determined experimentally. The differences average less than 0.1%, with one standard deviation of 5.3% when all experiments at ≥4.35 kbar are considered. Deviations are slightly skewed towards positive values. The standard deviation decreases to 4.9% when X(NaCl) ≤ 0.5 is considered, and to 3.9% when using only the 10 and 15 kbar data. Thus, Eqn. (3) is expected to provide an accurate value to within



Fig. 5. (A) Temperature-independent isobars of relative quartz solubility from Eqn. (3), text. (B) Difference, in percent, between relative silica molality predicted by Eqn. (3), text, and experimental determinations at the same temperatures and pressures. The low-pressure data (4.35 kbar) display increasing deviation from the equation with increasing  $X_{NaCl}$ . Although devations for the high-pressure (10–15 kbar) data are greatest at high  $X_{NaCl}$ , there is no systematic trend with fluid composition. The region between the dashed lines corresponds to two standard deviations in the difference for 10–15 kbar data.

about 10% at the  $2\sigma$  level at 500–900°C, 4–15 kb, and  $0 < X(\text{NaCl}) \le 0.5$ . It is not valid outside this range. Table 4 lists values of  $m(\text{SiO}_2)/m_o(\text{SiO}_2)$  calculated with Eqn. (3) at discrete pressure and X(NaCl) values.

Our analysis predicts that silica solubility at 4 kbar is nearly independent of salinity to X(NaCl) = 0.15 (Fig. 5A). A very

Table 3. Parameters for equation 1.

				1						
			i =							
		0	1	2	3	4				
	0 1	0.96907 10.418	0.012348 -4.0242	-0.00129 0.46876	3.65E-05 -0.02565	-7.31E-08 0.000539				
j =	2 3	-41.383 38.761	13.996 - 12.961	-1.689 1.6059	-0.095512 -0.09275	-0.00206 0.002024				

Table 4. Relative monanty (m <sub>P</sub> m <sub>o</sub> ) from equation 1.												
P (kbar)		X(NaCl)										
	0.05	0.10	0.15	0.20	0.25	0.30	0.35	0.40	0.45	0.50		
4	1.000	0.971	0.918	0.848	0.765	0.674	0.582	0.494	0.414	0.349		
5	0.954	0.892	0.821	0.744	0.663	0.581	0.502	0.428	0.361	0.305		
6	0.920	0.834	0.750	0.667	0.588	0.512	0.442	0.377	0.320	0.271		
7	0.895	0.792	0.698	0.612	0.533	0.461	0.397	0.340	0.289	0.246		
8	0.876	0.762	0.661	0.572	0.493	0.425	0.365	0.312	0.267	0.227		
9	0.862	0.740	0.634	0.543	0.465	0.398	0.341	0.293	0.251	0.214		
10	0.850	0.722	0.613	0.521	0.443	0.379	0.324	0.279	0.240	0.205		
11	0.840	0.708	0.596	0.504	0.427	0.364	0.312	0.269	0.232	0.199		
12	0.831	0.695	0.582	0.489	0.414	0.352	0.303	0.262	0.227	0.195		
13	0.823	0.684	0.570	0.478	0.403	0.344	0.296	0.257	0.223	0.191		
14	0.817	0.676	0.561	0.469	0.395	0.337	0.291	0.253	0.220	0.188		
15	0.813	0.671	0.556	0.464	0.391	0.334	0.288	0.251	0.218	0.186		

Table 4. Relative molality  $(m_i/m_o)$  from equation 1

small absolute salting-in may occur at low salinity at that pressure. Figure 5A shows that the salting-in tendency inherited from low pressures is expected to result in elevated relative solubility to no more than about 5-6 kbar, above which there

is a nearly exponential decrease in solubility with increasing X(NaCl). The loss of the salting-in effect with increasing pressure can also be visualized in a three-dimensional perspective view (Fig. 6).



Fig. 6. Perspective view of the temperature-independent relative quartz solubility surface as a function of  $X_{NaCl}$  and pressure.



Fig. 7. Silica molality as a function of NaCl mole fraction as predicted by Eqn. (3), text, at 500, 700, and 900°C and 4–15 kbar. The great regularity of solute behavior at  $X_{NaCl} > 0.2$  indicates that the fluids attain a simple constitution at high pressure and NaCl concentrations, which may be viewed as an ionic hydrosaline magma.

Figure 7 shows silica concentrations in H<sub>2</sub>O-NaCl fluids in equilibrium with quartz at 500, 700, and 900°C, 4–15 kbar, and X(NaCl) = 0–0.5 and Figure 8 shows SiO<sub>2</sub> molality isopleths at 500–900°C at 100°C intervals. Equation (3) was used in combination with predicted values of m<sub>o</sub>(SiO<sub>2</sub>) from Manning (1994) to generate these figures. Figures 7 and 8 illustrate that at 500–600°C, X(NaCl) > 0.2, there is very little difference in silica concentration over the entire pressure range at fixed X(NaCl). Thus, in general, pressure does not strongly influence quartz solubility in concentrated brines at low temperatures and P > 4 kb. At higher temperatures, the isopleths display progressively greater separation because m<sub>o</sub>(SiO<sub>2</sub>) increases with P and T (Fig. 8C–E).

#### 3.5. Solubility in CO<sub>2</sub>-H<sub>2</sub>O Solutions

Quartz solubility measurements in  $CO_2$ -H<sub>2</sub>O solutions at 10 kbar and 800°C (Table 2) are subject to larger uncertainties than the NaCl-H<sub>2</sub>O data at the same conditions because of the small solubilities, as well as the requirement of smaller crystals and fluid volumes. Trial extrapolation of the lower-pressure data of Shettel (1974) shows that, unlike in NaCl solutions at low pressues, solubility in the CO<sub>2</sub>-H<sub>2</sub>O system is fairly regular and predictable with changing T and P. The following equation, derived in part from the lower P,T data as summarized by Walther and Orville (1983), reproduces the present measurements in  $CO_2$ -H<sub>2</sub>O solutions at 10 kbar and 800°C within uncertainties:

$$m(SiO_2) = 0.0301 \{ \exp[3.76 \times (H_2O)] - 1 \}$$
(4)

Figure 9 shows  $m(SiO_2)$  versus  $H_2O$  activity,  $a(H_2O)$ . The standard state of  $H_2O$  is the pure phase at any temperature and pressure. Activities are consistent with Aranovich and Newton (1996) for NaCl solutions and Aranovich and Newton (1999) for CO<sub>2</sub> solutions. The large difference in silica solubility behavior between the two media occurs because of the opposite mixing properties of the CO<sub>2</sub>-H<sub>2</sub>O and NaCl-H<sub>2</sub>O systems at

elevated pressures: in NaCl solutions the  $H_2O$  activities are much lower than the mole fractions at pressures above 4 kbar, by virtue of pressure-induced ionic dissociation (Aranovich and Newton, 1996), whereas increasing pressure causes  $H_2O$  activities to exceed the mole fractions in  $CO_2$  solutions by substantial amounts above 5 kbar. Figure 9 shows that quartz is as much as an order of magnitude more soluble in concentrated NaCl solutions than in  $CO_2$  solutions of the same  $H_2O$  activity, in the low  $a(H_2O)$  range (<0.5).

#### 4. DISCUSSION AND INTERPRETATION

A principal finding of the present investigation is the change with pressure from initial enhancement of SiO<sub>2</sub> solubility in NaCl solutions at 2 kbar to initial salting-out at pressures above  $\sim$ 4 kbar, as shown in Figures 1 and 4. This behavior closely parallels the large decrease of H<sub>2</sub>O activity with increasing pressure, produced undoubtedly by dissociation of a neutral solute NaClo complex into Na<sup>+</sup> and Cl<sup>-</sup> ions as H<sub>2</sub>O densities increase above about 0.7 gm/cm3 (Aranovich and Newton, 1996). It seems highly probable, therefore, that the change from salting-in to salting-out is occasioned by dissociation of NaClo into ions. Indeed, two different solution regimes appear to exist, with contrasting effects on silica solubility. At low solution densities, encountered at low pressure and low salinity, the salting-in of SiO<sub>2</sub> is maximized. At high solution densities, change in silica solubility with increasing salinity is more regular, with monotonic salting-out. This fundamental change in solution behavior may be characterized as the transition from complex high-temperature aqueous solutions dominated by the neutral NaClo complex, to ionic solutions, or hydrosaline melts, exhibiting remarkably regular SiO<sub>2</sub> solubility behavior over large ranges of T, P, and salinity, as illustrated in Figures 3 and 7.

It is easily shown that the exponential decrease of quartz solubility with increasing NaCl concentration at 10 kbar is an expected effect if the solubility is very low in molten NaCl, and if the solution is a simple mixture of ions and molecules without excess enthalpy or non-configurational entropy. A significant change of slope of the plot of log SiO<sub>2</sub> molality versus  $H_2O$  mole fraction (Fig. 3), may exist in the 900°C isotherm below  $X(H_2O) = 0.5$ , possibly signalling a change in constitution of the solute silica species. An accelerated decrease in solubility at high salinity would imply that the dominant silica solute species is a hydrate, which is expected to become unstable at very low  $H_2O$  activities.

The nature of the silica solute species may be investigated in the plot of log  $m(SiO_2)$  versus log  $H_2O$  activity, following Walther and Orville (1983), as shown in Figure 10. The slope of the trend in  $H_2O$ -NaCl solutions suggests that the dominant solute species is  $H_4SiO_4$ , with possible admixture of a silica monohydrate at lower  $H_2O$  activities. The change to lower hydration state would be consonant with decreasing  $H_2O$  activity, which must progressively destabilize hydrous solute species. This analysis assumes that there is no significant Nabearing solute complex, such as that suggested by Anderson and Burnham (1983). Our data do not rule out the latter possibility.

A similar analysis for the present solubility data in  $CO_2$ -H<sub>2</sub>O solutions (Fig. 10) indicates that the dominant dissolved silica



Fig. 8. (A–E). Silica molality isopleths at 500–900°C as a function of pressure and  $X_{NaCl}$  as determined by Eqn. (3), text.

species is a neutral dihydrate with one to three attached (solvated)  $H_2O$  molecules, the solvation number decreasing with decreasing  $H_2O$  activity. This is virtually the same behavior as found by Walther and Orville (1983) at pressures of 5 kbar and below. It is apparent that the aqueous solubility behavior of

quartz is much different in NaCl-bearing than in  $\rm CO_2$ -bearing solutions.

The two disparate theories of  $SiO_2$  solubility in NaCl solutions may be contrasted, especially with regard to the origin of the salting-in effect. Xie and Walther (1993), after a detailed



Fig. 9. Comparison of silica solubility at 10 kbar and 800°C in concentrated NaCl-H<sub>2</sub>O and CO<sub>2</sub>-H<sub>2</sub>O solutions as functions of H<sub>2</sub>O activity ( $a_{H_2O}$ ). The plot shows that dissolved SiO<sub>2</sub> partitions preferentially into concentrated salt solutions relative to CO<sub>2</sub>-rich fluids. Because such fluids are immiscible over a wide range of high P-T, the data illustrate that silica will be strongly fractionated into the brine. The large error bars at low H<sub>2</sub>O activity result from the low solubility of quartz in CO<sub>2</sub>-rich fluids. H<sub>2</sub>O activities in NaCl solutions from Aranovich and Newton (1996) and in CO<sub>2</sub> solutions from Aranovich and Newton (1996).

analysis of the electrostatics of NaCl- $H_2O$  solutions, concluded that the neutral NaCl<sup>o</sup> complex does not have sufficient dipole moment to account for the magnitude of the observed salting-in effect at 0.5 to 2 kbar. They appealed instead to both solvation and short-range electrostatic interactions between dissociated



Fig. 10. Log m(SiO<sub>2</sub>) versus log a(H<sub>2</sub>O) at 800°C and 10 kbar. Present experimental data for CO<sub>2</sub>-H<sub>2</sub>O solutions lie in the sector between slopes of n = 3 and n = 5. Following Walther and Orville (1983), n is interpreted as the total hydration number of the dominant solute silica species, necessarily decreasing with increasing salinity. If the same interpretation is applicable to the data for NaCl solutions, the hydration number would decrease with decreasing H<sub>2</sub>O activity from 1.9 in pure H<sub>2</sub>O to 1.1 at a(H<sub>2</sub>O) = 0.3. The plot shows that an entirely different solution mode exists for quartz in concentrated electrolyte solutions at high pressures than for solutions in a solvent containing a neutral non-polar component.

 $Na^+$  and  $Cl^-$  ions and  $H_2O$  molecules, with resultant increase of the dielectric constant of the solution. In this approach, variations in silica molality with X(NaCl) are related through the Setchenow equation:

$$\log \gamma(\text{SiO}_2) = b' \text{ I} \tag{5}$$

In this equation, I is the ionic strength of the solution, defined as one-half the sum of the molalities of charged species, b' is the empirical Setchenow coefficient, which is a function of T and P, and adopting a standard state for aqueous SiO<sub>2</sub> of unit activity of the hypothetical one molal solution at infinite dilution at a given pressure and temperature,  $\gamma$ (SiO<sub>2</sub>) is the ratio of SiO<sub>2</sub> molality in pure H<sub>2</sub>O to that in H<sub>2</sub>O-NaCl solution. In Eqn. (5), negative values of b' promote small values of  $\gamma$ (SiO<sub>2</sub>) and hence salting-in.

Xie and Walther (1993) fitted Eqn. (5) to their solubility data at 0.5 and 1.0 kbar to explain the observed high-temperature salting-in and low-temperature salting-out. Their reconstruction assumes that no complexing takes place between dissolved silica and solute ions. The approach of Anderson and Burnham (1983) is quite different from that of Xie and Walther (1993) in that the salting-in would be explained by the SiO<sub>2</sub> solution reaction with NaCl, according to Eqn. (2).

Both theories have shortcomings in application to the present study. The Xie and Walther (1993) reconstruction must be regarded as mainly empirical, since the calculated Setchenow coefficients show large variation over relatively small ranges of T and P, and it is uncertain how to extrapolate the low pressure results to mid- and lower-crustal pressures. The solute reaction theory of Anderson and Burnham (1983) predicts that the solutions become quite acid for SiO<sub>2</sub>-rich solutions, which is questionable. An attempt by these authors to calculate an equilibrium constant for the reaction assumes that, apart from the salting-in effect produced by reaction (2), NaCl solutions would have the same effect on SiO<sub>2</sub> solubility as CO<sub>2</sub> solutions at the same X(H<sub>2</sub>O). Figure 10 shows that, for the conditions of 800°C and 10 kbar, where there is no salting-in, the solution mechanisms are quite different in the two media. Neither theory allows for the close correspondence between salting-out behavior and pressure-induced decrease of H<sub>2</sub>O activity.

In a more general way, it seems clear why addition of NaCl at high pressures should decrease the concentration of silica faster than at low pressures. In Eqn. (1), if n = 2, the equilibrium constant is  $K(T,P) = a(H_4SiO_4)/a^2(H_2O)$ . At ten kbar and high temperature, the H<sub>2</sub>O activity is nearly equal to the square of the mole fraction in NaCl solutions (Aranovich and Newton, 1996), whereas at 2 kbar, the activity is close to the mole fraction. Hence, the activity (and concentration) of H<sub>4</sub>SiO<sub>4</sub> will fall faster at 10 kbar with increasing salinity than at 2 kbar, with consequently greater tendency to overcome any salting-in effect.

#### 5. APPLICATION TO PETROLOGY

Several mechanisms have been suggested for generation of concentrated alkali-chloride solutions during crustal processes. At shallow depths, a hyperconcentrated brine or hydrosaline melt may result from fluid immiscibility at high temperatures, most plausibly in cooling of fluids of magmatic origin (Fournier, 1987). At deeper levels, pore fluids may be enriched in solutes by extraction of  $H_2O$  into anatectic melts (Philippot, 1993), by preferential absorption of  $H_2O$  in hydration reactions (Bennett and Backer, 1992; Markl and Bücher, 1998; Kullerud and Erambert, 1999), by prograde dehydration reactions in deeply subducted serpentinite and sea-floor sediments (Scambelluri et al., 1997), or as magmatic exhalations from crystallizing mafic magmas in deep-crustal environments (Hansteen and Burke, 1994). Although phase relations in the ternary system  $H_2O$ -NaCl-CO<sub>2</sub> are still imperfectly known at high temperatures and pressures, it is likely that fluid immiscibility exists between concentrated brines and CO<sub>2</sub>-rich fluids at most P-T conditions prevailing in high-grade metamorphism (Duan et al., 1995).

The present work shows that NaCl solutions at deep-crustal metamorphic conditions are substantial solvents for silica, despite the strong decrease of quartz solubility with NaCl concentration. At 10 kbar and 800°C, a concentrated brine of H<sub>2</sub>O mole fraction 0.6 will have a H<sub>2</sub>O activity of 0.45, according to Aranovich and Newton (1996), and, according to the experimental work of Aranovich and Newton (1998), the H<sub>2</sub>O activity is low enough to stabilize orthopyroxene in biotite-bearing rocks. This solution would dissolve about 0.7 wt.% of SiO<sub>2</sub> in quartz-bearing rocks, according to the present solubility study, or about the same as in pure water at 1 kbar and 800°C. Such fluids could account for the synmetamorphic orthopyroxenebearing quartz veins found in some classic granulite facies terranes (e.g., Bamble, S. Norway-Van den Kerkhof et al., 1994; S. India-Hansen et al., 1995). The alternative low-H<sub>2</sub>Oactivity fluid which has been often cited for granulite facies metamorphism, namely CO<sub>2</sub>-rich fluids, would be much more inert, with a silica solubility eight times lower at the same H<sub>2</sub>O activity. The vapor-absent interpretation of granulite facies metamorphism (Lamb and Valley, 1984) could not account for the quartz-orthopyroxene veins. The additional properties of high alkali mobility and grain wetting ability, again in contrast to CO<sub>2</sub>-rich fluids, make concentrated alkali chloride solutions seem feasible as chemically active agents in deep-crustal metamorphism.

An alkali-rich, low-H<sub>2</sub>O activity fluid provides a conceptual model for interpretation of many features of metamorphic rocks, including migmatites, albitization, myrmekite, and replacement antiperthite (Newton et al., 1998). The ability of hypersaline fluids to exchange alkalis in deep-seated environments without large-scale rock melting is enhanced by their extraordinary ability to infiltrate low-porosity silicate aggregates along grain boundaries. Metasomatic effects could be quite substantial, even for small fluid-rock ratios, because of the very high alkali contents of concentrated brines. The solubility of silica as feldspar-like molecules may be considerably higher than in the simple system  $SiO_2-H_2O-NaCl$ , and higher silica solubility is expected in alkaline solutions, such as those involved in fenitization. Silica solubility in more complex systems is an important area for further investigation.

The uniformity of silica solubility in concentrated NaCl solutions at 4 to 15 kbar is noteworthy. At 700°C, solutions of  $X(H_2O) = 0.6$  have a very small range of silica solubility at these pressures, less than 20 percent variation. The regularity of silica solubility in high-pressure brines is illustrated in Figures 3 and 7. Apparently, completely ionized brines are relatively simple substances, which become more complex in their be-

havior at lower pressures and fluid densities as the NaCl<sup>o</sup> neutral complex becomes stable. In this regard, a completely ionized NaCl solution may be a more convenient standard state for solute silica than the traditional dilute solution state at low pressure, which, as this study shows is not suitable for prediction of quartz solubility in brines over most of the pressure-temperature-concentration ranges of deep crust/upper mantle metamorphism.

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